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A LIGAND FIELD THEORY ANALYSIS OF THE SPECTRA OF THE T263 LEVEL--ETC(U)

MAR 78 J D WEBB, E R BERNSTEIN

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"A LIGAND FIELD THEORY ANALYSIS OF THE SPECTRA  
OF THE  $t_{2g}^3$  LEVELS OF  $IrF_6$ "

by

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five  $t_{2g}^3$  levels cannot be fit with physically reasonable Racah, spin orbit, and crystal field parameters. Parameter values presented are derived from a fit of the lowest three levels only:  $B = 297 \text{ cm}^{-1}$ ;  $C = 1167 \text{ cm}^{-1}$ ;  $\zeta_{5d} = 4182 \text{ cm}^{-1}$  and  $10 Dq = 35,000 \text{ cm}^{-1}$ . Gas-to-crystal shifts and site splittings are then calculated that further indicate strong influence of charge transfer states on the nature of the  $t_{2g}^3$  manifold. Finally, new data for the  $\Gamma_{8g} (^2E_g)$  and  $\Gamma_{6g} (^2T_{1g})$  states at  $1.2 \text{ } \mu\text{m}$  are given. Small Jahn-Teller interactions are observed in the  $\Gamma_{8g} (^2E_g)$  state at this wavelength and a substantial site splitting is assigned for it. Both of these effects further demonstrate the importance of charge transfer admixture in the  $t_{2g}$  manifold.

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## I. INTRODUCTION

$\text{IrF}_6$  has played an important role in the ligand field theory of  $5d^n$  transition metal compounds.<sup>1-8</sup> It is well established that the sharp, detailed spectra observed in the visible and near IR can be described approximately as transitions within a  $t_{2g}^3$  manifold.<sup>2,9</sup> However, a rigorous attempt to match observed and calculated<sup>6,10a,b</sup> electronic energies has not been made. Present availability of definitive assignments of electronic origins in both vapor<sup>9</sup> and mixed crystals of  $\text{IrF}_6/\text{MoF}_6$ <sup>11,12</sup> makes a detailed ligand field theory study feasible. Additional mixed crystal data for the  $\Gamma_{8g}$  ( $^2E_g$ ) and  $\Gamma_{6g}$  ( $^2T_{1g}$ ) levels of  $\text{IrF}_6$  are presented here to facilitate further comparison and complete the requisite data set.

There are several reasons why such a study might be useful. First, the most obvious benefit would be determination of reliable electrostatic (B, C) and spin-orbit ( $\zeta_{5d}$ ) parameters. Comparison with other members of the series ( $\text{ReF}_6$ ,  $\text{OsF}_6$ ,  $\text{PtF}_6$ ) might then allow insight into the nature of ligand-metal interactions in these high oxidation state (VI) compounds. Second, the extent to which charge-transfer (CT) states interact with  $t_{2g}^n$  ligand field states is still an open question. In the case of  $\text{CrBr}_3$  ( $3d^3$ ), it has been speculated that such an interaction accounts for the inability to fit observed  $t_{2g}^3$  levels with a ligand field calculation.<sup>13</sup> If interaction with CT states is important,  $\text{IrF}_6$  is a good candidate for study because an even CT state ( $\Gamma_{6g}$  or  $\Gamma_{7g}$ ) has recently been identified only  $\sim 1,000 \text{ cm}^{-1}$  above the  $\Gamma_{8g}$  ( $^2T_{2g}$ ) origin<sup>11</sup> at  $\sim 14,900 \text{ cm}^{-1}$  (the more intense odd CT bands begin at  $\sim 18,500 \text{ cm}^{-1}$ ). Third, since both vapor and mixed crystal data are available, gas-to-crystal shifts and site splittings can also be studied. These small electrostatic perturbations can be effectively employed to generate detailed information about the nature of the (nominally)  $t_{2g}^3$

levels of  $\text{IrF}_6$ . The magnitude, direction, and nature of the response of the electronic and vibronic states to crystal interactions will be a sensitive function of the actual state descriptions.

In Section IIIA, results of an attempt to fit observed  $t_{2g}^3$  levels of  $\text{IrF}_6$  with a parametric ligand field calculation are discussed. Gas-to-crystal shift data and low symmetry crystal field splitting of the  $\Gamma_{8g}$  levels are examined in Section IIIB and IIIC, respectively. Mixed crystal data for the  $\Gamma_{8g} (^2E_g)$  and  $\Gamma_{6g} (^2T_{1g})$  levels of  $\text{IrF}_6/\text{MoF}_6$  are discussed in detail in Section IIID.

## II. METHODS AND RESULTS

### A. Experimental

A summary of data and assignments for the  $\Gamma_{8g}$  ( $^2E_g$ ) and  $\Gamma_{6g}$  ( $^2T_{1g}$ ) levels of  $\text{IrF}_6/\text{MoF}_6$  are given in Table 1. This table completes the  $\text{IrF}_6/\text{MoF}_6$  mixed crystal data for  $t_{2g}^3$  levels of  $\text{IrF}_6$ .<sup>11,12</sup> These spectra are also presented in Figures 1 and 2. Experimental methods employed are identical to those presented in reference 11.

### B. Calculational

Results of an attempt to fit observed  $t_{2g}^3$  ligand field levels of  $\text{IrF}_6$  with a parametric calculation are given in Tables 2 and 3. The energy levels were calculated using Eisenstein's<sup>10a,b</sup>  $d^3$  energy matrices; however, due to the large ligand field splitting ( $10 Dq \sim 35,000 \text{ cm}^{-1}$ ), only one set of Racah and spin-orbit parameters ( $B, C, \zeta_{5d}$ ) were used. A linear least squares routine was employed in parameter variation for an optimum fit. The results given in Table 2 were obtained by optimizing the fit of all five excited  $t_{2g}^3$  levels while those in Table 3 were obtained by optimizing the fit of only the lowest three levels.

An attempt to calculate gas-to-crystal shifts for  $\text{IrF}_6$  and  $\text{ReF}_6$  (Tables 4 and 5) by varying the ligand field parameter  $Dq$  is summarized in Table 6. The  $\text{IrF}_6$   $B, C$ , and  $\zeta_{5d}$  parameters employed were those obtained for the calculation reported in Table 3. A spin-orbit parameter value of  $3,200 \text{ cm}^{-1}$  was used for the  $\text{ReF}_6$  gas-to-crystal shift calculations. A tetragonal crystal field calculation for a  $d^3$  system was also carried out in order to evaluate site splitting mechanisms. Since appropriate matrix elements for a strong field-double group basis set are not available, these are given in Table 7. Tetragonal field matrix elements were generated by the appropriate basis transformation of Rahman's<sup>14</sup>  $d^3$  tetragonal matrix elements for a strong field-

single group basis. Transformed tetragonal matrix elements may then be combined with Runciman and Schroeder's<sup>6</sup> octahedral  $d^3$  matrices, since Tanabe and Sugano's phase is used throughout.<sup>15</sup> An example of this calculation is given in Table 8. The appropriate one-electron tetragonal parameters  $\rho$  and  $\mu$  are as follows:

$$\left. \begin{aligned} \langle t_{2g\zeta} | \hat{O}(e_g, \theta) | t_{2g\zeta} \rangle &= -\frac{2}{3} \rho \\ \langle t_{2g\xi} | \hat{O}(e_g, \theta) | t_{2g\xi} \rangle &= \langle t_{2g\eta} | \hat{O}(e_g, \theta) | t_{2g\eta} \rangle = \frac{1}{3} \rho \end{aligned} \right\} \quad (1)$$

$$\langle e_{g\theta} | \hat{O}(e_g, \theta) | e_{g\theta} \rangle = - \langle e_{g\epsilon} | \hat{O}(e_g, \theta) | e_{g\epsilon} \rangle = -\frac{1}{2} \mu \quad (2)$$



### III. DISCUSSION

#### A. Calculations of $t_{2g}^3$ electronic levels of $\text{IrF}_6$

Results of  $\text{IrF}_6$  calculations (Table 2) indicate several difficulties with the parametric fit of all the  $t_{2g}^3$  levels. The main problem seems to be with the Racah parameters (e.g.,  $B(\text{IrF}_6) < 0$ ), although the low spin-orbit value ( $\zeta_{5d}(\text{IrF}_6) = 3012 \text{ cm}^{-1}$ ) compared with that of  $\text{ReF}_6^4$  ( $\zeta_{5d} \sim 3200 \text{ cm}^{-1}$ ) and the large average absolute standard deviation ( $\sigma \sim 250 \text{ cm}^{-1}$ ) are also disturbing. Note, however, that values obtained for the Racah parameters are not quite as poor as might appear at first glance since only  $3B + C$  is determined directly within the  $t_{2g}^3$  block. The unphysical negative value for  $B$  and the resulting high value for  $C$  do, however, indicate that the  $t_{2g}^3$  levels are perturbed.

A possible cause of the above difficulties might be thought to be neglect of the Trees' correction, which is important for 3d transition metal ions<sup>13</sup>. The calculation for  $\text{ReF}_6^{-2}$  in Table 2, however, demonstrates that such is not the case for these 5d systems.

The most plausible explanation for the above perturbation is a substantial interaction between the charge-transfer (CT) and ligand field electronic states. If this is the case, one expects that ligand field calculations for  $\text{ReF}_6^{-2}$   $t_{2g}^3$  states will be more accurate than those for  $\text{ReCl}_6^{-2}$  and  $\text{IrF}_6$  (see Table 2) due to the proximity of the CT and upper  $t_{2g}^3$  states in the latter two systems (see Table 9). Since only the mixing of even CT states into the  $5d-t_{2g}^3$  manifold is expected to be important in this context, this correlation requires that a rough equality exists between the onset energies of even and odd CT states in the two ionic systems. This is found to be the case for  $\text{IrF}_6$ .<sup>11</sup>

Another calculation can be done which indicates more directly that CT states do indeed perturb and mix with the  $t_{2g}^3$  manifold. Ligand field para-

meters in this calculation are optimized to fit only the lowest three excited levels (Table 3), as it is expected that these levels are less perturbed by CT bands. This type of calculation for  $\text{IrF}_6$  (Table 3) supports the CT hypothesis in that the resulting parameters ( $\zeta_{5d} = 4000 \text{ cm}^{-1}$ ,  $\frac{C}{B} = 4$ ), are more reasonable than those obtained for the fit of all five  $\text{IrF}_6$  levels ( $\zeta_{5d} = 3000 \text{ cm}^{-1}$ ,  $\frac{C}{B} = 15$ ). Of course, calculated energies for the  $\Gamma_{7g} (^2T_{2g})$  and  $\Gamma_{8g} (^2T_{2g})$  levels are much higher than those observed. The analogous calculations for  $\text{ReF}_6^{-2}$  and  $\text{ReCl}_6^{-2}$  yield results expected in this theory: the  $\Gamma_7 (^2T_{2g})$  and  $\Gamma_8 (^2T_{2g})$  calculated levels for  $\text{ReCl}_6^{-2}$  are far from the observed, while those for  $\text{ReF}_6^{-2}$  are substantially closer.

These considerations all lead to the conclusion that there is a substantial interaction between even CT and  $t_{2g}^3$  ligand field electronic levels of an isolated gas phase ( $0_h^*$ )  $\text{IrF}_6$  molecule. It is expected that similar interactions exist in other hexahalides with low-lying CT bands (i.e.,  $\text{ReCl}_6^{-2}$ ,  $\text{ReBr}_6^{-2}$ ,  $\text{PtF}_6$ , and  $\text{OsF}_6$ ).

In view of the magnitude of the shifts of the  $\Gamma_{7g} (^2T_{2g})$  and  $\Gamma_{8g} (^2T_{2g})$  levels of  $\text{IrF}_6$ , it is reasonable to assume that the lower levels are also significantly shifted, and thus that the parameter values obtained for  $\text{IrF}_6$  in Table 3 are only approximate. It may be said, however, that the present spin-orbit parameter value ( $\zeta_{5d} = 4182 \text{ cm}^{-1}$ ) is more realistic than that derived by Jorgensen<sup>8</sup> ( $\zeta_{5d} = 3100 \text{ cm}^{-1}$ ) who based his estimate on the energies of the  $\Gamma_{7g} (^2T_{2g})$  and  $\Gamma_{8g} (^2T_{2g})$  levels of  $\text{IrF}_6$ .

#### B. Gas-to-crystal shifts of the $t_{2g}^3$ levels of $\text{IrF}_6$

The gas-to-crystal (GC) shift data for mixed crystals of  $\text{IrF}_6$  in  $\text{UF}_6$ ,  $\text{WF}_6$ , and  $\text{MoF}_6$  are presented in Tables 4a and 5a; analogous data for  $\text{ReF}_6$  are given in Tables 4b and 5b. The  $\text{IrF}_6$  mixed crystal data have two notable features:

GC shifts of the lowest three electronic levels are virtually equal in each host (Table 4a), and the normalized (to the smallest shift in each host) sets of shifts are nearly equal in the three hosts (see Table 5). The shifts themselves appear to be linearly dependent on a single quantity which is varying from host to host. Note that the  $\text{ReF}_6$  GC shifts are nearly two orders of magnitude smaller than those found for  $\text{IrF}_6$ .

The crystal field (to be distinguished from ligand field) which acts on an  $\text{IrF}_6$  molecule at a crystal site is a reasonable candidate for the cause of GC shifts. Thus, the form of the crystal field operator and its effect on various approximate  $\text{IrF}_6$  wavefunctions will be examined. Site symmetry of all hexafluoride crystals (except 295 K neat  $\text{IrF}_6$ ) employed in this study is  $C_s (\sigma_d)$ .<sup>11</sup> Since  $\text{IrF}_6$  is octahedral in the gas phase, the crystal field operator which acts on the  $t_{2g}$  electrons of  $\text{IrF}_6$  may be expressed in terms of  $O_h^*$  tensor operators<sup>16</sup> as follows:

$$\begin{aligned} \hat{O}(\text{CF}) = & \hat{O}(\text{octahedral}) + \hat{O}(\text{tetragonal}) + \hat{O}(\text{trigonal}) + \hat{O}'(t_{2g}, \zeta) + \\ & \{\hat{O}(t_{1g}, x) - \hat{O}(t_{1g}, y)\} + \hat{O}(\text{ungerade terms}) \end{aligned} \quad (3)$$

in which

$$\hat{O}(\text{octahedral}) = \hat{O}(a_{1g})$$

$$\hat{O}(\text{tetragonal}) = \hat{O}(e_g, \theta)$$

$$\hat{O}(\text{trigonal}) = \{\hat{O}(t_{2g}, \xi) + \hat{O}(t_{2g}, \eta) + \hat{O}(t_{2g}, \zeta)\}$$

and  $\hat{O}'(t_{2g}, \zeta)$  is needed to maintain rigorous  $C_s^*$  site symmetry. Gerade terms are given in detail since they are typically the more important; ungerade terms are usually important only in intensity considerations.<sup>16</sup>

The 295 K neat  $\text{IrF}_6$  crystal GC shift data (Table 5) provide evidence necessary for the determination of the relative importance of each of the terms in Equation 3. Above -273 K all second and third row transition metal hexa-



fluoride crystals exhibit a body-centered cubic modification.<sup>17</sup> More detailed crystal structure data on  $\text{MoF}_6$ <sup>18</sup> imply that in the cubic modification, hexafluoride molecules are undistorted from an octahedral configuration. This argues that non-totally symmetric terms in the crystal field are not of major importance in the cubic crystal. Nonetheless the GC shifts observed in 295 K neat  $\text{IrF}_6$  crystals are substantial, about 50% of the shift observed at 2 K.

The conclusion which may be drawn from the above physical arguments is that the totally-symmetric crystal field operator  $\hat{O}(a_{1g})$  is probably the dominant one for these considerations. Moreover, it may be shown that none of the non-totally symmetric terms may lead to diagonal matrix elements which could cause an overall GC shift. Non-totally symmetric operators will thereby cause a GC shift only through off-diagonal matrix elements involving sizeable energy denominators. Consequently, subsequent discussion will consider only the effect of  $\hat{O}(a_{1g})$ .

The GC shifts calculated with the  $\hat{O}(a_{1g})$  operator depend critically on the functions chosen to represent the ligand field states. The first approximate wavefunctions of  $\text{IrF}_6$  which will be considered are those involving only  $t_{2g}^3$  functions (i.e., all configuration interaction is neglected). If the secular matrix for the octahedral crystal field operator is formed with these functions, it is found to be diagonal with each diagonal element equal to  $-4Dq'$ ; thus, at this level of approximation no GC shift occurs.

If complete ligand field wavefunctions are considered (i.e.,  $e_g$  configuration interaction included) then, as illustrated in Table 6, GC shifts are predicted. However, this mechanism can be rejected for three reasons: the increase in the octahedral field necessary to produce shifts of the magnitude



observed is larger than what is reasonable ( $10 Dq' \sim 15,000 \text{ cm}^{-1}$ ); calculated shifts do not match the observed GC shifts well; and concomitant GC shifts in  $\text{ReF}_6$  are not observed.

Considerations in Section IIA led to the conclusion that in gas phase  $\text{IrF}_6$ ,  $t_{2g}^3$  levels are perturbed by nearby even charge-transfer (CT) bands; CT character is thereby mixed into these ligand field free molecule wavefunctions. The final approximate wavefunctions considered are these ligand field-CT functions. The GC shift data may be explained if sufficient CT character is present in the above functions to cause diagonal  $\hat{O}(a_{1g})$  matrix elements to deviate substantially from the  $-4 Dq'$  value mentioned above. However, since there is not an adequate theory which would allow a quantitative assessment of this mechanism, only qualitative remarks are possible. Three points favoring this interpretation are:

- 1) The highest two  $t_{2g}^3$  levels of  $\text{IrF}_6$  have been strongly repelled by CT bands ( $\sim 1900$  and  $\sim 600 \text{ cm}^{-1}$ , Table 3). Presumably, the lower three levels are also repelled, both by CT bands and by perturbed  $\Gamma_{7g} (^2T_{2g})$  and  $\Gamma_{8g} (^2T_{2g})$  levels. The magnitudes of level repulsions imply substantial mixing of  $t_{2g}^3$  and CT wavefunctions.
- 2) The magnitude of the GC shifts is largest for the  $\Gamma_{8g} (^2T_{2g})$  and  $\Gamma_{7g} (^2T_{2g})$  levels which are energetically closest to the CT bands.
- 3) The linear behavior of GC shifts with host-to-host variation noted above (see Table 5) agrees with the simple energy expressions implied by this mechanism. These energy expressions depend linearly on an octahedral crystal field parameter which is expected to evidence small variations as a function of host.

On the other hand, the proposed CT- $t_{2g}^n$  mixing mechanism for GC shifts does not appear to be supported by  $\text{ReF}_6$  GC shift data (Table 5b). The  $\text{ReF}_6$  GC shifts are not as large as might be expected based on standard second-order perturbation theory and comparison with the GC shift for the  $\text{IrF}_6$   $t_{2g}^3 e_g^2$  ( ${}^2T_{1g}$ ) level considering the energies of the respective CT bands (Table 9). It is however, difficult to comment on this apparent lack of agreement; there are too many unknown quantities involved. Little is known about the detailed nature of the interaction between CT and ligand field states and how the interaction might change from  $\text{IrF}_6$  to  $\text{ReF}_6$ . Thus, in spite of this point, the CT-admixture mechanism constitutes a reasonable interpretation of the GC shift data.

C. Low symmetry crystal field splitting of the  $t_{2g}^3 e_g^2$  levels of  $\text{IrF}_6$

Site splittings of the electronic origins of  $t_{2g}^3 e_g^2$  levels of  $\text{IrF}_6$  in various host crystals are given in Table 10. In contrast to the situation outlined in Section IIIB, non-totally symmetric crystal field operators in Equation 3 are important for this effect. Calculations of site splitting associated with a general low symmetry crystal field (Equation 3) are prohibitively difficult because of the large number of parameters involved. However, it has been demonstrated by crystal structure ( $\text{UF}_6$ )<sup>19</sup> and spectroscopic ( $\text{MoF}_6$ )<sup>11</sup> data that the ( $e_g, \theta$ ) tetragonal component of the field is probably a reasonable approximation to the field experienced by  $\text{IrF}_6$  at host sites. The tetragonal crystal field matrix elements for a complete ligand field basis given in Table 7 were employed in this analysis. Results of this calculation are given in Table 8 and may be compared with experimental data in Table 10. The observed and calculated splitting of the  $t_{2g}^3 e_g^2$  ( ${}^2E_g$ ) level are found to be not even in qualitative agreement. In light of preceding discussion, this discrepancy may be interpreted in terms of CT-admixture

into the  $t_{2g}^3$  levels. Clearly both  $t_{2g}^3$  and charge-transfer states contribute to the observed splittings throughout the manifold. A detailed decomposition of the effects of CT-admixture, tetragonal crystal field, and their interplay with the Jahn-Teller interaction is not apparent at present. However, this subject is treated to some extent in reference 20.

D. Absorption spectra of the  $\Gamma_{8g}$  ( $^2E_g$ ) and  $\Gamma_{6g}$  ( $^2T_{1g}$ ) levels of  $\text{IrF}_6/\text{MoF}_6$

Existence of a substantial site splitting of the  $\Gamma_{8g}$  ( $^2E_g$ ) level at  $8200 \text{ cm}^{-1}$  was important for the assessment of crystal field charge-transfer effects on  $\text{IrF}_6$   $\Gamma_{8g}$  states (Section IIIC). Since assignment of a  $30.6 \text{ cm}^{-1}$  site splitting of this level is not entirely obvious from the origin spectrum presented in Figure 1, it is necessary to indicate how this determination is made. Such an assignment becomes evident from an examination of the bending region (Figure 2 and Table 1) in which  $\nu_6$  and  $\nu_4$  are seen to have equally intense components separated by roughly  $30 \text{ cm}^{-1}$ . It should be noted that observation of crystal field splitting of the level at  $\sim 8200 \text{ cm}^{-1}$  provides the first strong evidence that the assignment of the level as  $\Gamma_{8g}$  ( $^2E_g$ ) is correct. Previous arguments were based on the observation of one component of the  $\nu_2$  ( $e_g$ ) vibration and the supposition that its change in energy from that observed in the ground state is due to a Jahn-Teller interaction.<sup>9</sup>

The linewidths, especially of the origins and  $\nu_6$ , are larger than is typical for the rest of the  $t_{2g}^3$  manifold. This may be due to the proximity of  $\Gamma_{8g}$  ( $^2E_g$ ) and  $\Gamma_{6g}$  ( $^2T_{1g}$ ) levels to the  $\Gamma_{8g}$  ( $^2T_{1g}$ ) level at  $1.6 \mu\text{m}$ . Vibronic coupling between these states could lead to fast relaxation which would broaden the transitions.



Broad lines and overlap of the  $\Gamma_{8g} (^2E_g)$  and  $\Gamma_{6g} (^2T_{1g})$  manifolds have inhibited interpretation of the vibronic portion of the spectrum. Nonetheless, a weak Jahn-Teller interaction ( $D_5 = 0.03$ )<sup>11</sup> may be identified for the  $\nu_5 (t_{2g})$  vibration in the  $\Gamma_{8g} (^2E_g)$  electronic manifold (see Figure 2 and Table 1).



#### IV. CONCLUSIONS

The main conclusion which may be drawn from this investigation is that the interaction of even charge-transfer (CT) bands with the  $t_{2g}^3$  levels of free molecule  $\text{IrF}_6$  is sufficiently strong to effect their detailed behavior. The following appear to be manifestations of the admixture of CT character into the  $t_{2g}^3$  levels:

- 1) Failure of ligand field theory to give a reasonable parametric fit to all of the observed gas-phase  $t_{2g}^3$  energies.
- 2) Substantial gas-to-crystal shifts for the  $t_{2g}^3$  levels.
- 3) Tetragonal crystal field splitting of the  $\Gamma_{8g} (^2E_g)$  level.

In spite of the CT interaction, rough ligand field theory parameters (B, C,  $\zeta_{5d}$ ) have been determined based on the lower three excited  $t_{2g}^3$  levels (Table 3). It is found that the spin-orbit parameter is larger than previously believed.

Conclusive evidence for the assignment of the  $\Gamma_{8g} (^2E_g)$  level of  $\text{IrF}_6$  is found, based on site splitting of the origin. A small linear JT interaction has been identified in this state for the  $\nu_5$  mode ( $D_5 \sim 0.03$ ).

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Table 1. The absorption spectra of the  $\Gamma_{8g} (^2E_g)$  and  $\Gamma_{6g} (^2T_{1g})$  levels of  $\text{IrF}_6/\text{MoF}_6$  at 1.6 K with 1  $\text{cm}^{-1}$  slitwidths. The accuracy is approximately  $\pm 0.5 \text{ cm}^{-1}$ .

$\lambda_{\text{Air}} (\text{\AA})$	$\sigma_{\text{Vacuum}} (\text{cm}^{-1})$	I (a)	FWHH (b)	$\Delta\sigma_1 (\text{cm}^{-1})$ (c)	$\Delta\sigma_2 (\text{cm}^{-1})$ (d)	Assignments
12218.9	8181.8	M	10	0		Origin(a) $\{\Gamma_{8g} (^2E_g)\}$
12173.3	8212.4	W	7	30.6		Origin(b) $\{\Gamma_{8g} (^2E_g)\}$
11908.2	8395.3	M	16	213.5		$\nu_6$ (a) $\{\Gamma_{8g} (^2E_g)\}$
11869.5	8422.6	-	18	240.8		$\nu_5(J_5=3/2)\{\Gamma_{8g} (^2E_g)\}$
11862.5	8427.6	M		245.8		$\nu_6$ (b) $\{\Gamma_{8g} (^2E_g)\}$
11814.5	8461.9	W	16	280.1		$\nu_5(J_5=1/2)\{\Gamma_{8g} (^2E_g)\}$
11808.9	8465.9	M		284.1		$\nu_4$ (a) $\{\Gamma_{8g} (^2E_g)\}$
11767.3	8495.8	M	13	314.0		$\nu_4$ (b) $\{\Gamma_{8g} (^2E_g)\}$
11460.9	8722.9	M	12	541.1	0	Origin $\{\Gamma_{6g} (^2T_{1g})\}$
11394.4	8773.8	W		592.1	51.0	Phonons $\{\Gamma_{6g} (^2T_{1g})\}$
11295.7	8850.5	W	7	668.7	127.6	$\nu_2\{\Gamma_{8g} (^2E_g)\} (?)$
11206.5	8920.9	W		739.2	198.1	
11182.7	8939.9	S	21	758.1	217.0	$\nu_6\{\Gamma_{6g} (^2T_{1g})\}$
11164.5	8954.5	M		772.7	231.6	
11121.3	8989.3	M		807.5	266.4	$\nu_5\{\Gamma_{6g} (^2T_{1g})\}$
11105.8	9001.8	S		820.0	278.9	$\nu_4\{\Gamma_{6g} (^2T_{1g})\}$
11081.3	9021.7	W		839.9	298.8	
11039.3	9056.1	W		874.3	333.2	$(\nu_2 + \nu_6)(a)\{\Gamma_{8g} (^2E_g)\} (?)$

continued...

Table 1. (continued)

$\lambda_{\text{Air}}(\text{\AA})$	$\sigma_{\text{Vacuum}}(\text{cm}^{-1})$	I (a)	FWHH (b)	$\Delta\sigma_1(\text{cm}^{-1})$ (c)	$\Delta\sigma_2(\text{cm}^{-1})$ (d)	Assignments
10957.4	9123.8	W		941.9	400.8	$(\nu_2 + \nu_4)(a) \{ \Gamma_{8g} (^2E_g) \} (?)$
10859.3	9206.2	W		1024.4	483.3	$(\nu_5 + \nu_6) \{ \Gamma_{6g} (^2T_{1g}) \}$
10789.9	9265.4	W		1083.6	542.5	$(\nu_5 + \nu_4) \{ \Gamma_{6g} (^2T_{1g}) \}$
10612.1	9420.6	W		1238.8	697.7	$\nu_3 \{ \Gamma_{6g} (^2T_{1g}) \}$

a) Intensity: S = strong, M = medium, W = weak.

b) FWHH = full width at half-height.

c) Energy ( $\text{cm}^{-1}$ ) relative to the  $\Gamma_{8g} (^2E_g)$  origin (a).d) Energy ( $\text{cm}^{-1}$ ) relative to the  $\Gamma_{6g} (^2T_{1g})$  origin.



Table 2. Comparison of ligand field calculations of the  $t_{2g}^3$  levels of some  $5d^3$  systems ( $\text{IrF}_6$ ,  $\text{ReF}_6^{-2}$ ,  $\text{ReCl}_6^{-2}$ ). Eisenstein's<sup>(a)</sup> energy matrix for  $d^3$  was used along with a linear least squares fitting routine which optimized the fit for all five levels. B and C are Racah's electrostatic parameters,  $\zeta_{5d}$  is the spin-orbit coupling parameter, and  $10 Dq$  is the ligand field parameter. Only B, C,  $\zeta_{5d}$  were varied. The various  $\Gamma_{ig}$  label the  $t_{2g}^3$  levels. All the values are given in wavenumbers.

a) $\text{IrF}_6$ (Vapor)	$\Gamma_{8g}(^2T_{1g})$	$\Gamma_{8g}(^2E_g)$	$\Gamma_{6g}(^2T_{1g})$	$\Gamma_{7g}(^2T_{2g})$	$\Gamma_{8g}(^2T_{2g})$
Observed: <sup>(b)</sup>	6261	8333	8858	12328	15156
Calculated:	6561	8397	8180	12345	15344

Parameters:  $B = -221$ ,  $C = 3188$ ,  $\zeta_{5d} = 3012$ ,  $10 Dq = 35000$ <sup>(c)</sup>,  $3B + C = 2525$ .  
Avg. Abs. Deviation<sup>(g)</sup>: 250.

b) $\text{ReF}_6^{-2}$					
Observed:	9080 <sup>(b)</sup>	10130 <sup>(b)</sup>	11160 <sup>(b)</sup>	17390 <sup>(d)</sup>	18670 <sup>(d)</sup>
Calculated:	9033	10278	11197	17329	18641

Parameters:  $B = 665$ ,  $C = 1504$ ,  $\zeta_{5d} = 3019$ ,  $10 Dq = 33100$ <sup>(e)</sup>,  $3B + C = 3499$ .  
Avg. Abs. Deviation: 64.

c) $\text{ReCl}_6^{-2}$					
Observed: <sup>(f)</sup>	7600 <sup>(b)</sup>	8906	9344	13840	15298
Calculated:	7748	8839	8844	13952	15443

Parameters:  $B = 4$ ,  $C = 2820$ ,  $\zeta_{5d} = 2254$ ,  $10 Dq = 30347$ <sup>(f)</sup>,  $3B + C = 2832$ .  
Avg. Abs. Deviation: 194.

a) References 10a, b.

b) Reference 9.

c) The  $10 Dq$  value of  $35000 \text{ cm}^{-1}$  was chosen on the basis of  $\text{ReF}_6^{-2}$  data<sup>(d)</sup>,  $\text{ReF}_6$  data<sup>(2)</sup>, and expected trends. The results are not very sensitive to the value of this parameter.

d) J. A. LoMenzo, S. Strohbridge, H. H. Patterson, J. Mol. Spec. 66, 150 (1977).

e) The  $10 Dq$  value used is taken from ref. d. The parameter values for B, C,  $\zeta$  found here are probably less realistic than those in ref. d since their data on the  $^4T_{2g}$  levels was utilized. However, this calculation is for comparison purposes with the  $\text{IrF}_6$  calculation only.

f) P. B. Dorain, R. G. Wheeler, J. Chem. Phys. 45, 1172 (1966).

g) Average Absolute Deviation of the calculated energies from the observed energies.

Table 3. Comparison of the ligand field calculations of the  $t_{2g}^3$  levels of some  $5d^3$  systems ( $\text{IrF}_6$ ,  $\text{ReF}_6^{-2}$ ,  $\text{ReCl}_6^{-2}$ ). This table is similar to Table 2 except that the linear least square fit was optimized for only the lowest three  $t_{2g}^3$  levels. The predicted energies of the highest two levels are given in parentheses. Note that the highest two levels are calculated as being much higher than observed for  $\text{IrF}_6$  and to a lesser extent for  $\text{ReCl}_6^{-2}$  and  $\text{ReF}_6^{-2}$  (see text).

a) $\text{IrF}_6$		$r_{8g}(^2T_{1g})$	$r_{8g}(^2E_g)$	$r_{6g}(^2T_{1g})$	$r_{7g}(^2T_{2g})$	$r_{8g}(^2T_{2g})$
(a)						
Observed:		6261	8333	8858	12328	15156
Calculated:		6275	8322	8793	(12922)	(17084)

Parameters:  $B = 297$ ,  $C = 1167$ ,  $\zeta_{5d} = 4182$ ,  $10 Dq = 35000$ ,  $3B + C = 2058$ .

b) $\text{ReF}_6^{-2}$						
Observed:		9080 <sup>(b)</sup>	10130 <sup>(b)</sup>	11160 <sup>(b)</sup>	17390 <sup>(c)</sup>	18670 <sup>(c)</sup>
Calculated:		9080	10130	11160	(17102)	(18046)

Parameters:  $B = 792$ ,  $C = 1179$ ,  $\zeta_{5d} = 2745$ ,  $10 Dq = 33000$ ,  $3B + C = 3555$ .

c) $\text{ReCl}_6^{-2}$						
(d)						
Observed:		7600 <sup>(b)</sup>	8906	9344	13840	15298
Calculated:		7600	8906	9344	(14704)	(16349)

Parameters:  $B = 333$ ,  $C = 1873$ ,  $\zeta_{5d} = 2792$ ,  $10 Dq = 31000$ ,  $3B + C = 2872$ .

a) Reference 9.

b) Reference 8.

c) Reference d, Table 2.

d) Reference f, Table 2.

Table 4. a) Energies of the  $t_{2g}^3$  levels of  $\text{IrF}_6$  in the gas phase and in various crystals. For the  $\Gamma_{8g}$  levels, which are slightly split by a low-symmetry crystal field, the center of gravity is given.

b) Energies of the  $t_{2g}$  levels of  $\text{ReF}_6$  in the gas phase and in various crystals. The split  $\Gamma_{8g}$  levels are treated as above.

a) $\text{IrF}_6$	$\Gamma_{8g} ({}^4A_2)$	$\Gamma_{8g} ({}^2T_{1g})$	$\Gamma_{8g} ({}^2E_g)$	$\Gamma_{6g} ({}^2T_{1g})$	$\Gamma_{7g} ({}^2T_{2g})$	$\Gamma_{8g} ({}^2T_{2g})$
$\text{IrF}_6$ (Vapor) <sup>(a)</sup>	0	6261	8333	8858	12328	15156
Neat $\text{IrF}_6$ <sup>(a)</sup> (-2.2 K)	0	6114	8177	8701	12060	14878
Neat $\text{IrF}_6$ <sup>(b)</sup> (295 K)	0	6188	8256	8779	12177	14947
$\text{IrF}_6/\text{UF}_6$ <sup>(b)</sup>	0	6111	8185	8708	12082	14883
$\text{IrF}_6/\text{WF}_6$ <sup>(b)</sup>	0	6135	8206	8730	12118	14926
$\text{IrF}_6/\text{MoF}_6$	0	6123 <sup>(c)</sup>	8194 <sup>(d)</sup>	8720 <sup>(d)</sup>	12093 <sup>(e)</sup>	14901 <sup>(e)</sup>

b) $\text{ReF}_6$	$\Gamma_{8g} ({}^2T_{2g})$	$\Gamma_{7g} ({}^2T_{2g})$
$\text{ReF}_6$ (Vapor) <sup>(f)</sup>	0	5001
$\text{ReF}_6/\text{UF}_6$ <sup>(g)</sup>	0	5003
$\text{ReF}_6/\text{WF}_6$ <sup>(g)</sup>	0	5001
$\text{ReF}_6/\text{MoF}_6$ <sup>(g)</sup>	0	4997

a) Reference 9.

b) E. R. Bernstein, J. D. Webb, unpublished results.

c) Reference 11.

d) Table 1.

e) Reference 12.

f) J.C.D. Brand, G. L. Goodman, B. Weinstock, J. Mol. Spec. 38, 449 (1971).

g) E. R. Bernstein, G. R. Meredith, J. Chem. Phys. 64, 375 (1976).

Table 5. a) Gas-to-crystal shifts in  $\text{cm}^{-1}$  ( $\Delta\Gamma = \Gamma(\text{crystal}) - \Gamma(\text{gas})$ ) for mixed and neat crystals of  $\text{IrF}_6$  (see Table 4a). The parenthetical numbers represent the shifts as normalized to the smallest shift for a given host.

a)  $\text{IrF}_6$

	$\Delta\Gamma_{8g}(^2T_{1g})$	$\Delta\Gamma_{8g}(^2E_g)$	$\Delta\Gamma_{8g}(^2T_{1g})$	$\Delta\Gamma_{7g}(^2T_{2g})$	$\Delta\Gamma_{8g}(^2T_{2g})$
MIXED CRYSTALS					
$\text{IrF}_6/\text{UF}_6$	-150 (1.01)	-148 (1.00)	-150 (1.01)	-246 (1.66)	-272 (1.84)
$\text{IrF}_6/\text{WF}_6$	-126 (1.00)	-127 (1.01)	-128 (1.02)	-210 (1.67)	-230 (1.82)
$\text{IrF}_6/\text{MoF}_6$	-138 (1.00)	-139 (1.01)	-138 (1.00)	-235 (1.70)	-255 (1.85)
NEAT CRYSTALS					
$\text{IrF}_6(\sim 2.2 \text{ K})$	-147 (1.00)	-156 (1.06)	-157 (1.07)	-268 (1.82)	-278 (1.89)
$\text{IrF}_6(295 \text{ K})$	-73 (1.00)	-77 (1.05)	-79 (1.08)	-151 (2.07)	-209 (2.86)

b) Gas-to-crystal shifts for mixed crystals of  $\text{ReF}_6$  (see Table 4b).

	$\Delta\Gamma_{7g}(^2T_{2g})$
$\text{ReF}_6/\text{UF}_6$	2
$\text{ReF}_6/\text{WF}_6$	0
$\text{ReF}_6/\text{MoF}_6$	3



Table 6. Calculated energies ( $\text{cm}^{-1}$ ) of the upper  $t_{2g}^3$  levels of  $\text{IrF}_6$  and the higher  $t_{2g}$  level of  $\text{ReF}_6$  as a function of  $10 Dq$ . Full  $5d^3 e_g-t_{2g}$  wave functions and configuration interaction are used. The electrostatic and spin-orbit parameters (all in wavenumbers) used for  $\text{IrF}_6$  are the following:  $B = 297$ ,  $C = 1167$ ,  $\zeta_{5d} = 4182$ . The spin-orbit parameter used for  $\text{ReF}_6$  is  $\zeta_{5d} = 3200$ . The parenthetical numbers represent a gas-to-crystal shift if  $10 Dq = 35,000 \text{ cm}^{-1}$  represents the gas phase value, while the  $10 Dq$  values across the top of the table are taken to represent those in the crystal.

$10 Dq$	25000	30000	35000	40000	45000	50000
<b><math>\text{IrF}_6</math></b>						
$\Gamma_{8g}({}^2T_{1g})$	6390 (115)	6324 (49)	6275 (0)	6238 (-37)	6209 (-66)	6185 (-90)
$\Gamma_{8g}({}^2E_g)$	8438 (116)	8370 (48)	8322 (0)	8285 (-37)	8275 (-47)	8234 (-88)
$\Gamma_{6g}({}^2T_{1g})$	9028 (235)	8895 (102)	8793 (0)	8712 (-81)	8646 (-147)	8592 (-201)
$\Gamma_{7g}({}^2T_{2g})$	13096 (174)	13006 (84)	12922 (0)	12849 (-73)	12786 (-136)	12733 (-189)
$\Gamma_{8g}({}^2T_{2g})$	17341 (257)	17205 (121)	17084 (0)	16983 (-101)	16899 (-185)	16829 (-255)
<b><math>\text{ReF}_6</math></b>						
$\Gamma_{7g}({}^2T_{1g})$	5365 (150)	5278 (63)	5215 (0)	5166 (-49)	5127 (-88)	5096 (-119)

Table 7. Tetragonal  $d^3$  crystal field matrix elements in a strong field-double group  $O_h^*$  basis. The matrix elements were obtained by transforming Rahman's (a) tetragonal matrix elements from a single group to a double group basis.  $\rho$  and  $\mu$  are defined by Rahman (a) and in the text. The ordering of the basis functions is that of Runciman and Schroeder (b) since these matrix elements are combined with theirs to obtain the full  $d^3$  matrices. The  $\Gamma_6$  ( $O_h^*$ ) and  $\Gamma_7$  ( $O_h^*$ ) basis functions are numbered from 22-30, while  $\Gamma_8$  ( $O_h^*$ ) are numbered from 1-21. The  $D_{4h}^*$  matrix elements are obtained in 2 blocks,  $\Gamma_{6g}$  ( $D_{4h}^*$ ) and  $\Gamma_{7g}$  ( $D_{4h}^*$ ). The correlation between  $O_h^*$  and  $D_{4h}^*$  is such that:  $\Gamma_{6g} \rightarrow \Gamma_{6g}$ ,  $\Gamma_{7g} \rightarrow \Gamma_{7g}$ , and  $\Gamma_{8g} \rightarrow \Gamma_{6g} + \Gamma_{7g}$ .

$\Gamma_{6g}$ ( $D_{4h}^*$ ): $\Gamma_{8g}$ ( $O_h^*$ ) - $\Gamma_{8g}$ ( $O_h^*$ ) Sub-block			
$E6(3,4) = -\frac{\sqrt{3}}{3} \rho$	$E6(6,8) = -\frac{1}{5} \mu$	$E6(12,12) = \frac{1}{3} \rho + \frac{1}{4} \mu$	$E6(17,17) = -\frac{1}{3} \rho$
$E6(5,5) = -\frac{4}{15} \rho + \frac{1}{5} \mu$	$E6(7,7) = \frac{4}{15} \rho + \frac{1}{5} \mu$	$E6(12,14) = \frac{1}{4} \mu$	$E6(18,18) = -\frac{1}{3} \rho$
$E6(5,6) = \frac{1}{5} \rho - \frac{3}{20} \mu$	$E6(7,8) = -\frac{1}{5} \rho - \frac{3}{20} \mu$	$E6(13,13) = -\frac{1}{3} \rho - \frac{1}{4} \mu$	$E6(18,19) = \frac{1}{2} \mu$
$E6(5,7) = \frac{1}{5} \mu$	$E6(8,8) = -\frac{4}{15} \rho - \frac{1}{5} \mu$	$E6(14,14) = -\frac{1}{3} \rho + \frac{1}{4} \mu$	$E6(19,19) = \frac{1}{3} \rho$
$E6(5,8) = -\frac{3}{20} \mu$	$E6(9,10) = \frac{2}{3} \rho$	$E6(15,15) = \frac{4}{15} \rho$	$E6(19,20) = -\frac{1}{2} \mu$
$E6(6,6) = \frac{4}{15} \rho - \frac{1}{5} \mu$	$E6(11,11) = \frac{1}{3} \rho - \frac{1}{4} \mu$	$E6(15,16) = -\frac{1}{5} \rho$	$E6(20,20) = \frac{1}{3} \rho$
$E6(6,7) = -\frac{3}{20} \mu$	$E6(11,13) = -\frac{1}{4} \mu$	$E6(16,16) = -\frac{4}{15} \rho$	$E6(21,21) = \frac{1}{2} \mu$

Table 7. (continued)

$\Gamma_{6g} (D_{4h}^*): \Gamma_{8g} (0_h^*) - \Gamma_{6g} (0_h^*)$ Sub-block			
$E6(4,22) = \sqrt{\frac{6}{3}} \rho$	$E6(7,23) = -\sqrt{\frac{5}{20}} \mu$	$E6(10,25) = \frac{2}{3} \rho + \frac{1}{2} \mu$	$E6(15,28) = \sqrt{\frac{5}{15}} \rho$
$E6(5,23) = -\sqrt{\frac{5}{15}} \rho + \sqrt{\frac{5}{20}} \mu$	$E6(7,24) = \sqrt{\frac{5}{5}} \rho + \frac{3\sqrt{5}}{20} \mu$	$E6(11,26) = \sqrt{\frac{2}{3}} \rho - \sqrt{\frac{2}{4}} \mu$	$E6(16,28) = \sqrt{\frac{5}{5}} \rho$
$E6(5,24) = -\frac{3\sqrt{5}}{20} \mu$	$E6(8,23) = -\frac{3\sqrt{5}}{20} \mu$	$E6(12,27) = \sqrt{\frac{2}{3}} \rho + \sqrt{\frac{2}{4}} \mu$	$E6(17,29) = -\sqrt{\frac{2}{3}} \rho$
$E6(6,23) = -\sqrt{\frac{5}{5}} \rho + \frac{3\sqrt{5}}{20} \mu$	$E6(8,24) = -\sqrt{\frac{5}{15}} \rho - \sqrt{\frac{5}{20}} \mu$	$E6(13,26) = \sqrt{\frac{2}{4}} \mu$	$E6(18,30) = -\sqrt{\frac{2}{3}} \rho$
$E6(6,24) = \sqrt{\frac{5}{20}} \mu$	$E6(9,25) = -\frac{2}{3} \rho$	$E6(14,27) = -\sqrt{\frac{2}{4}} \mu$	$E6(19,30) = -\sqrt{\frac{2}{2}} \mu$

$\Gamma_{7g} (D_{4h}^*): \Gamma_{8g} (0_h^*) - \Gamma_{8g} (0_h^*)$  Sub-block

$E7(i,j) = -E6(i,j) \quad i, j \leq 21$

Table 7. (continued)

$\Gamma_{7g} (D_{4h}^*): \Gamma_{8g} (0^*) - \Gamma_{7g} (0^*)$ Sub-block			
$E7(3,22) = -\frac{\sqrt{5}}{3} \rho$	$E7(7,24) = -\frac{\sqrt{5}}{15} \rho - \frac{\sqrt{5}}{20} \mu$	$E7(12,27) = \frac{\sqrt{2}}{4} \mu$	$E7(19,29) = -\frac{\sqrt{2}}{3} \rho$
$E7(5,23) = -\frac{\sqrt{5}}{5} \rho + \frac{3\sqrt{5}}{20} \mu$	$E7(8,23) = \frac{\sqrt{5}}{20} \mu$	$E7(13,26) = \frac{\sqrt{2}}{3} \rho + \frac{\sqrt{2}}{4} \mu$	$E7(19,30) = \frac{\sqrt{2}}{2}$
$E7(5,24) = \frac{\sqrt{5}}{20} \mu$	$E7(8,24) = -\frac{\sqrt{5}}{5} \rho - \frac{3\sqrt{5}}{20} \mu$	$E7(14,27) = \frac{\sqrt{2}}{3} \rho - \frac{\sqrt{2}}{4} \mu$	$E7(20,29) = \frac{\sqrt{2}}{2} \mu$
$E7(6,23) = \frac{\sqrt{5}}{15} \rho - \frac{\sqrt{5}}{20} \mu$	$E7(9,25) = -\frac{2}{3} \rho$	$E7(15,28) = \frac{\sqrt{5}}{5} \rho$	$E7(20,30) = -\frac{\sqrt{2}}{3} \rho$
$E7(6,24) = \frac{3\sqrt{5}}{20} \mu$	$E7(10,25) = \frac{2}{3} \rho + \frac{1}{2} \mu$	$E7(16,28) = -\frac{\sqrt{5}}{15} \rho$	
$E7(7,23) = -\frac{3\sqrt{5}}{20} \mu$	$E7(11,26) = -\frac{\sqrt{2}}{4} \mu$	$E7(18,29) = \frac{\sqrt{2}}{2} \mu$	

a) Reference 14.

b) Reference 6.



Table 8. Site splitting of  $\Gamma_{8g}$  levels of  $\text{IrF}_6$  - results of a tetragonal crystal field calculation for the  $t_{2g}^3$  levels with the parameters:  $B = 340$ ,  $C = 1117$ ,  $\zeta_5 = 3987$ ,  $10 Dq = 35000 \text{ cm}^{-1}$  (a), and tetragonal parameters  $\rho$  and  $\mu$ . These latter parameters are defined in Equations 1 and 2 and are given here in  $\text{cm}^{-1}$  (see text). The splittings of the  $\Gamma_8$  states as a function of these parameters is given in  $\text{cm}^{-1}$ .

$\rho$	$\mu$	$\Gamma_8 (^4A_2)$	$\Gamma_8 (^2T_{1g})$	$\Gamma_{8g} (^2E_g)$	$\Gamma_8 (^2T_{2g})$
52	0	10	28	0.1	38
52	100	8.7	27	0.1	38
52	500	2.6	25	0.1	37

- a) These parameter values were chosen to match the  $\text{IrF}_6/\text{MoF}_6$  data with a Jahn-Teller correction of  $100 \text{ cm}^{-1}$  added to the  $\Gamma_{8g} (^2T_{1g})$  experimental energy. However, the results, especially the splitting of the  $\Gamma_{8g} (^2E_g)$  level, are not sensitive to the choice of parameters.

Table 9. Onset frequencies ( $\text{cm}^{-1}$ ) of the intense charge-transfer bands in various hexahalides.

	$\text{IrF}_6$ (a)	$\text{ReF}_6$ (a)	$\text{ReF}_6^{-2}$ (b)	$\text{ReCl}_6^{-2}$ (b)	$\text{UF}_6$ (c)	$\text{WF}_6$ (d)	$\text{MoF}_6$ (d)
$\sigma_{\text{CT}}$	18500	24000	>35000	21000	24500	60000	50000

a) Reference 2.

b) Reference 8. It is stated in reference e that the near UV bands in  $\text{ReCl}_6^{-2}$  are ligand field bands rather than charge transfer bands; however, recent work(f) has shown that the original charge transfer assignment(8) is correct.

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Table 10. Splittings of  $\Gamma_{8g}$  levels of  $\text{IrF}_6$  in various mixed crystals by a low symmetry crystal field, given in wavenumbers. Compare with Table 8.

	$\Gamma_{8g}({}^4A_2)$	$\Gamma_{8g}({}^2T_{1g})$	$\Gamma_{8g}({}^2E_g)$	$\Gamma_{8g}({}^2T_{2g})$
$\text{IrF}_6/\text{UF}_6$ <sup>(a)</sup>	10.0	61.9	53.0	66.4
$\text{IrF}_6/\text{WF}_6$ <sup>(a)</sup>	5.7	38.6	34.8	46.3
$\text{IrF}_6/\text{MoF}_6$	5.2 <sup>(b)</sup>	34.7 <sup>(c)</sup>	30.6 <sup>(d)</sup>	42.1 <sup>(b)</sup>

a) E. R. Bernstein and J. D. Webb, unpublished results.

b) Reference 12.

c) Reference 12.

d) Table 1.

Figure 1.

Origin of the  $\Gamma_{8g}$  ( $^2E_g$ ) electronic state of  $\text{IrF}_6/\text{MoF}_6$ . The origin is split by a low symmetry crystal field; see Figure 2 for verification of this assignment.



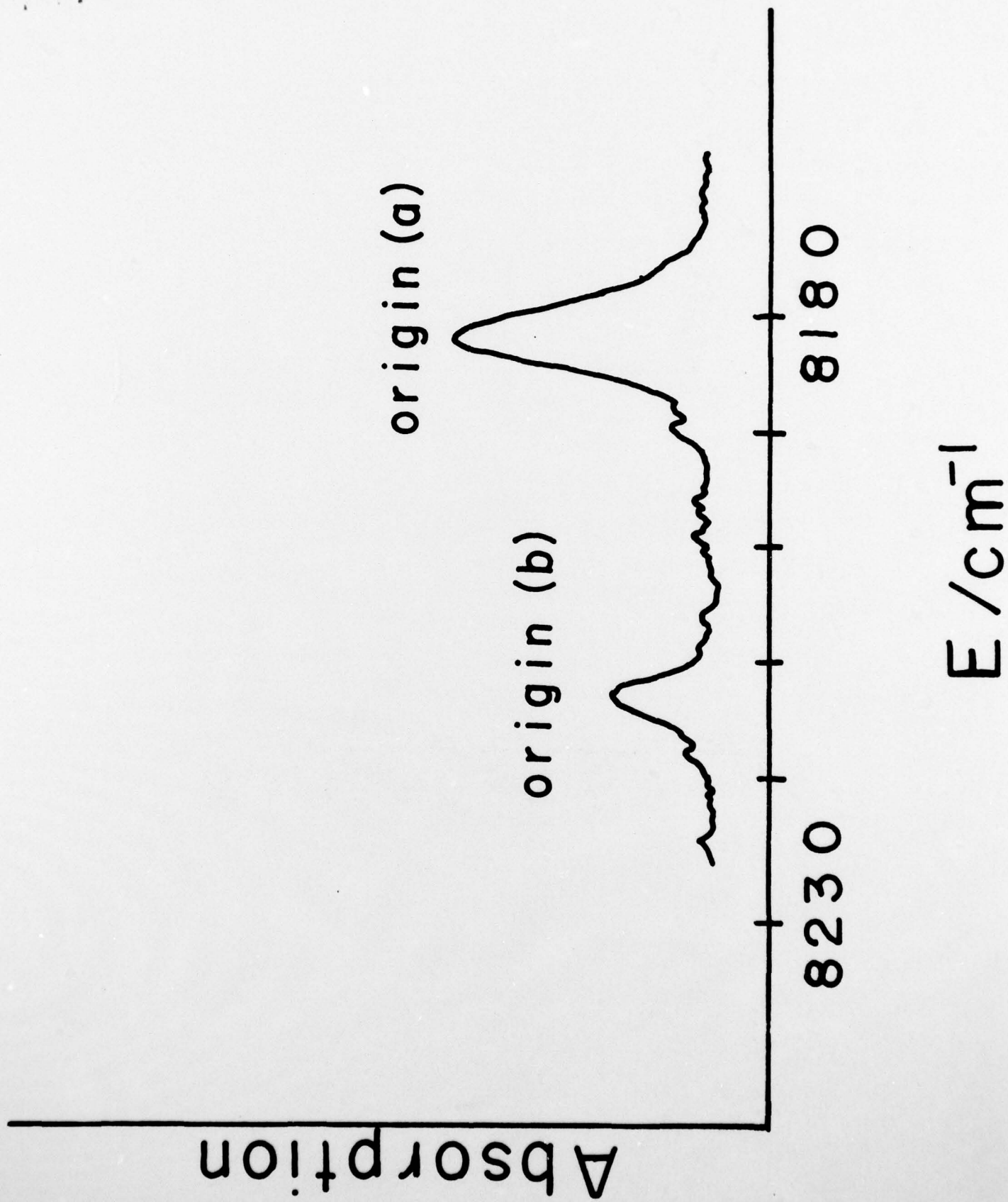
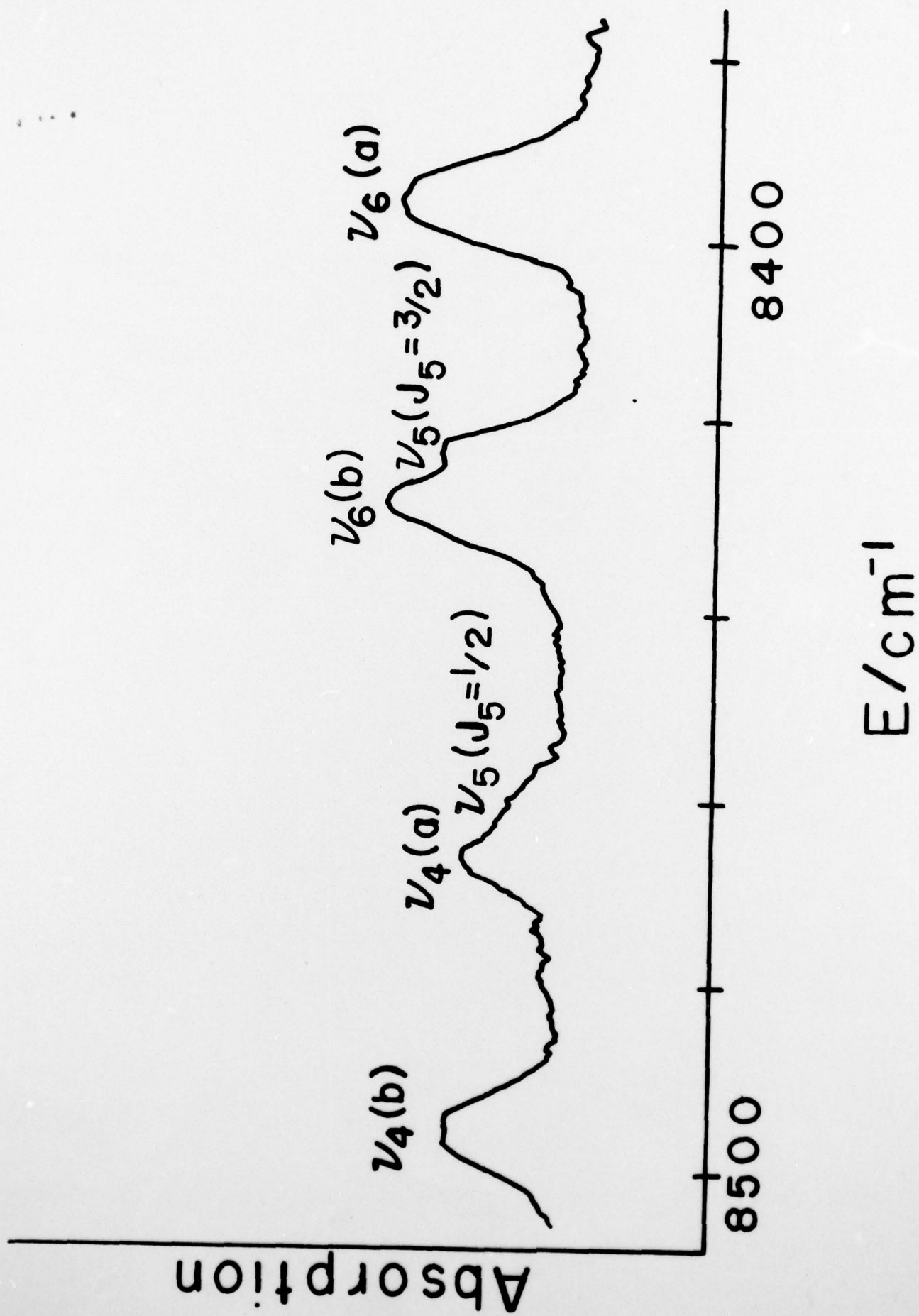


Figure 2.

Vibrational bending region of the  $r_{8g}$  ( $^2E_g$ ) electronic state of  $IrF_6/MoF_6$ . Note that the  $\sim 30\text{ cm}^{-1}$  spacing of the  $\nu_6(t_{2u})$  and  $\nu_4(t_{1u})$  components matches the origin splitting, verifying that the observed lines are due to a low symmetry crystal field splitting. A small linear Jahn-Teller splitting of  $\nu_5(t_{2g})$  is also apparent.



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